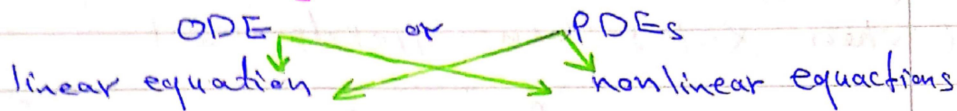


Lecture 6:

Self-Similar Scaling Solutions of DEs



separation of variables

similarity method

Laplace transformation

traveling wave soln

fourier transformation

For nonlinear problems, we will show that if the PDE can be scaled in such a way as to exactly reproduce its original form (as suggested by the term 'self-similar'), then it will possess a class of solutions that share this property. A consequence of the property of self-similarity is that these solutions can be obtained from a reduced version of the original problem. For a PDE with two independent variables, this leads to an ODE problem for the similarity solution. Despite their nature as exact solutions to a given problem under special conditions, they often provide important understanding of the broader behaviour of all solutions of the system.

The process of constructing a similarity solution to a given problem consists of three stages:

- (i) Looking for a scaling "symmetry" of the problem to see if similarity solutions are possible.
- (ii) Determining the forms of the similarity variables and functions from the scale-invariant IT groups for the problem.
- (iii) Solving a reduced problem for the similarity function and then transforming back to give the solution of original problem.

Finding scaling - Invariant Symmetries

For $u = u(x, t)$, will admit a similarity solution by making use of the change of variables

$$u(x, t) = U \tilde{u}(\tilde{x}, \tilde{t}), \quad x = L \tilde{x}, \quad t = T \tilde{t} \quad \text{--- (1)}$$

where U, L, T are undetermined real positive parameters. We call the problem scale invariant if relationships exist between the scaling parameters U, L, T in (1) that make the scaled problem take exactly the same as the original problem with at least one scaling parameter remaining undetermined.

problem for $u(x,t)$ = problem for $\tilde{u}(\tilde{x}, \tilde{t})$ --- (2)

For equation equations, it expected that the time scale T is the free parameter and the other scales can be expressed and $L = L(T)$.

Example: $u_t + u u_x = 0$ $0 \leq x < \infty$ --- (3)

$$u(0, t) = 0, \quad \int_0^{\infty} u(x, t) dx = 4 \quad \text{--- (4)}$$

Sol By using (1), we have

$$\tilde{u}_{\tilde{t}} + \left(\frac{uT}{L}\right) \tilde{u} \tilde{u}_{\tilde{x}} = 0 \quad 0 \leq \tilde{x} < \infty \quad \text{--- (5)}$$

$$\tilde{u}(0, \tilde{t}) = 0, \quad (uL) \int_0^{\infty} \tilde{u}(\tilde{x}, \tilde{t}) d\tilde{x} = 4 \quad \text{--- (6)}$$

setting $\frac{uT}{L} = 1$ and $uL = 1$

element the scaling and make the scaled system identical with (5) and (6) and hence the problem scale invariant.

$$\text{If } b=0 \Rightarrow \pi_2 = UT^c \Rightarrow \{\pi_6\} = UT^c = T^c$$

$$\rightarrow [\pi_2] = T^{-\frac{1}{2}+c} \rightarrow \boxed{c = \frac{1}{2}}$$

$$\text{If } f = \pi_2 \Rightarrow \left(f = t^{-\frac{1}{2}} u \right)$$

$$\text{From above } F(\pi_1, \pi_2) = 0$$

Applying the implicit function theorem gives

$$f = f(\eta) \rightarrow t^{-\frac{1}{2}} u = f(x, t^{-\frac{1}{2}})$$

$$\rightarrow u(x, t) = t^{\frac{1}{2}} f(x, t^{\frac{1}{2}}) \quad \text{--- (3)}$$

If we select $c=0 \rightarrow ??$

$$u(x, t) = ? \quad \text{H.W}$$

Example / $u_t = u_{xx}$ use the scaling (1) to find η and solve if with

$$\textcircled{1} \int_{-\infty}^{\infty} u(x, t) dx = \int_{-\infty}^{\infty} u(x, 0) dx = M$$

$$\textcircled{2} u(0, t) = 1 \quad u \rightarrow 0 \quad \text{as } x \rightarrow \infty$$

$$\rightarrow L = T^{\frac{1}{2}} \quad \rightarrow V = T^{-\frac{1}{2}}$$

$$\rightarrow U(x, t) = T^{-\frac{1}{2}} \tilde{U}(T^{-\frac{1}{2}}x, T^{-\frac{1}{2}}t)$$

for arbitrary of T . This transformation is called scaling symmetry of (3).

We define a scale-invariant parameter Π as a minimal product powers of the variables of the system

$$\Pi = U^a x^b t^c \quad \text{--- (8a)}$$

Let $[\Pi]$ represent the dependence of Π on the scaling

$$[\Pi] = U^a L^b T^c \quad \text{--- (8b)}$$

In particular for (7), this yields

$$[\Pi] = T^{-\frac{a}{2}} T^{\frac{b}{2}} T^c = T^0$$

If $a=0, b=1 \xrightarrow{\text{from (8)}} \Pi = x t^c \Rightarrow [\Pi] = L T^c$

$$[\Pi_1] = T^{\frac{1}{2}+c} \rightarrow \boxed{c = -\frac{1}{2}}$$

If $\eta = \Pi_1 \Rightarrow \boxed{\eta = x t^{-\frac{1}{2}}}$

$$\text{If } b=0 \Rightarrow \pi_2 = UT^c \Rightarrow \{\pi_0\} = UT^0 = T^0$$

$$\rightarrow [\pi_2] = T^{-\frac{1}{2}+c} \rightarrow \boxed{c = \frac{1}{2}}$$

$$\text{If } f = \pi_2 \Rightarrow \left(f = t^{-\frac{1}{2}} u \right)$$

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$$\rightarrow u(x, t) = t^{\frac{1}{2}} f(x, t^{-\frac{1}{2}}) \quad \text{--- (2)}$$

If we select $c=0 \rightarrow ??$

$$u(x, t) = ? \quad \text{H.W}$$

Example / $u_t = u_{xx}$ use the scaling (1) to find η and solve it with

$$\textcircled{1} \int_{-\infty}^{\infty} u(x, t) dx = \int_{-\infty}^{\infty} u(x, 0) dx = M \quad \text{--- (1)}$$

$$\textcircled{2} u(0, t) = 1 \quad u \rightarrow 0 \quad \text{as } x \rightarrow \infty$$