Chapter4 -General Motion of Particle in Three Dimensions

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4.1 Introduction: General Principles

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4.2 The Potential Energy Function in Three Dimensions Motion: The Del Operator

4.1 Introduction: General Principles:

 We now study the general case of motion of particle in three dimensions. The vector form of the equation of motion (Newton's Second Law)for such particle is :

$$F = \frac{dp}{dt} \qquad \qquad \mathbf{OR} \quad F = ma \tag{4.1}$$

as p = mv, is the linear momentum of the particle. In Cartesian coordinates, we can write:

 $F_x = m\ddot{x}$ $F_y = m\ddot{y}$ $F_z = m\ddot{z}$ (4.2)

The simplest solution is assuming a function (F) of spatial coordinates only(i.e. F=F(r)).

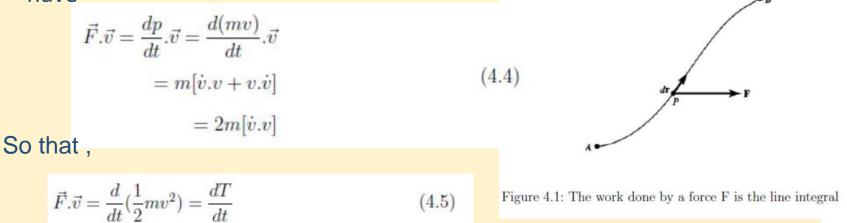
 Z_p

 y_p

4.1.1 The work Principle:

 $\vec{F} \cdot \frac{d\vec{r}}{dt} = \frac{dT}{dt}$

Work done on a particle causes it to gain or lose kinetic energy. Now, we will use equation (1) with taking the dot product with velocity (v), so, we have



Where **T** is the Kinetic Energy. However, v = dr/dt, so equation (4.5) becomes

(4.6)
$$\int \vec{F} \cdot d\vec{r} = \int dT \qquad (4.7)$$
$$W = T_f - T_i \qquad (4.8)$$

4.1.2 Conservative Force and Force Fields:

• If the forces acting on a particle were conservative , it could be drive from the derivative of a scalar potential function as following:

$$F_x = -\frac{dV(x)}{dx}$$
(4.9)
$$W = \int F_x dx = -\int dV(x)$$
$$W = \int F_x dx = -\Delta V(x) = V(A) - V(B)$$
(4.10)

 This mean, we need to know the potential energy at started point (A) and the ended point (B). From equation (8) and (10), we can find a general form of the conservation law for the total energy of the particle:

$$\int \vec{F} \cdot d\vec{r} = \int dT \qquad (4.7)$$

$$W = T_f - T_i \qquad (4.8)$$

$$W = \int F_x dx = -\bigtriangleup V(x) = V(A) - V(B) \qquad (4.10)$$

$$E_{total} = V(A) + T(A) = V(B) + T(B) = \text{Constant at the particle motion}$$

$$(4.11)$$

4.2 The Potential Energy Function in Three Dimensions Motion: The Del Operator

 Assume we have F and V in three dimensions so we can write:

As F(r)=iFx + jFy + kFz $F(r) = -\frac{dV(r)}{dr}$ (4.12), V(r)=iVx + jVy + kVz $-\frac{\partial V(x)}{\partial x}$ (4.13) $i \frac{\partial V(x)}{\partial x} - j \frac{\partial V(y)}{\partial y} -$ (4.14)

Del Operator

$$\nabla = i\frac{\partial}{\partial x} + j\frac{\partial}{\partial y} + k\frac{\partial}{\partial z}$$
(4.15)

Del Operator Properties

- $1 \nabla F =$ is a numerical quantity
- $2 \nabla \times F =$ **Curl F** ; lead to a new Vector
- $3 \nabla \times (\nabla F) = 0$; Curl of any grediat =0

(4.16)

Del Operator

Now, we express <u>Curl F</u> using Del operator

$$\mathbf{Curl} \ \mathbf{F} = \nabla \times F = \begin{vmatrix} i & j & k \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ F_x & F_y & F_z \end{vmatrix}$$
(4.17)

$$\nabla \times F = i(\frac{\partial F_z}{\partial y} - \frac{\partial F_y}{\partial z}) + j(\frac{\partial F_x}{\partial z} - \frac{\partial F_z}{\partial x}) + k(\frac{\partial F_y}{\partial x} - \frac{\partial F_x}{\partial y}) = 0$$
(4.18)

 Now, we can generalize the conservation of energy principle to three dimensions. The work done by a conservative force in moving a particle from point A to point B can be written as :

$$\int_{A}^{B} \vec{F}.dr = -\int_{A}^{B} \nabla V(\vec{r}).dr \quad \text{as} \quad F = -\nabla V$$

$$= \int_{A}^{B} (i\frac{\partial V}{\partial x} + j\frac{\partial V}{\partial y} + k\frac{\partial V}{\partial z}) \cdot (idx + jdy + kdz)$$
(4.19)

• H.W

$$\int_{A}^{B} dV(\vec{r}) = \Delta V]_{A}^{B} = V(B) - V(A)$$
(4.20)

• The last step illustrates the fact that $\nabla V. dr$ is an exact deferential equal to dV. The work done by any net force is always equal to the change in kinetic energy

i.e.
$$\int_{A}^{B} F.dr = \triangle T = -\triangle V$$

$$(4.21)$$

$$\triangle (T+V) = 0$$

T(A) + V(A) = T(B) + V(B) = E = Constant(4.22)

 We have arrived at our desired law of conservation of total energy. NOTE: If F is nonconservative force, it can not be equal to (−∇V) this lead to the work (F. dr) is not an exact deferential and can not be equaled to (dV). So, the general form of work -energy theorem becomes:

$$\int_{A}^{B} F.dr = \triangle(T+V) = -\triangle E \tag{4.23}$$

Example (1)

Given the two-dimensional potential energy function

$$V(\mathbf{r}) = V_0 - \frac{1}{2}k\delta^2 e^{-r^2/\delta^2}$$

where $\mathbf{r} = \mathbf{i} x + \mathbf{j} y$ and V_0 , k, and δ are constants, find the force function.

Solution:

We first write the potential energy function as a function of x and y,

$$V(x,y) = V_0 - \frac{1}{2}k\delta^2 e^{-(x^2 + y^2)/\delta^2}$$

and then apply the gradient operator:

$$F = -\nabla V$$

$$F = -\left(i\frac{\partial V}{\partial x} + j\frac{\partial V}{\partial y}\right)$$

$$V(x,y) = V_0 - \frac{1}{2}k\delta^2 e^{-(x^2+y^2)/\delta^2}$$

$$F = -\left(i\frac{\partial V}{\partial x} + j\frac{\partial V}{\partial y}\right)$$

$$\frac{\partial V}{\partial x} = \frac{1}{2}k\delta^2\left(\frac{2x}{\delta^2}\right)e^{-(x^2+y^2)/\delta^2}$$

$$\frac{\partial V}{\partial y} = \frac{1}{2}k\delta^2\left(\frac{2y}{\delta^2}\right)e^{-(x^2+y^2)/\delta^2}$$

$$\frac{\partial V}{\partial y} = ky e^{-(x^2+y^2)/\delta^2}$$

$$F = -(ikx \ e^{-(x^2 + y^2)/\delta^2} + jky \ e^{-(x^2 + y^2)/\delta^2})$$

 $F = -k(ix + jy)e^{-(x^2+y^2)/\delta^2}$

