

## Problems and Exercises Ch.4

**Problem 1:** Evaluate (a)  $\langle x \rangle$ , (b)  $\langle p_x^2 \rangle$ , (c)  $\langle p_x \rangle$  and (d)  $\langle p_x^2 \rangle$  for the eigenstates of a harmonic oscillator.

**Solution:** (a)  $\langle x \rangle = \int_{-\infty}^{\infty} \psi_n^*(x) x \psi_n(x) dx \quad \dots(6.37)$

We know that the harmonic oscillator wave functions have definite parity. Thus  $\psi_n(x)$  is either an even or an odd function of  $x$ . Therefore, the product  $\psi_n^*(x) \psi_n(x)$  will always be even. Since  $x$  is odd, the integrand will always be odd and hence  $\langle x \rangle = 0$ .

For a complete solution from first principles, we proceed as under.

We shall use the recurrence relation (6.23b) for the Hermite polynomials which can be rewritten as

$$2\alpha x H_n(\alpha x) = H_{n+1}(\alpha x) + 2n H_{n-1}(\alpha x) \quad \dots(6.38)$$

The harmonic oscillator wave functions can be written in terms of the Hermite polynomials as

$$\psi_n(x) = \left( \frac{\alpha}{2^n n! \sqrt{\pi}} \right)^{1/2} e^{-\alpha^2 x^2 / 2} H_n(\alpha x), \quad n = 0, 1, 2, \dots \quad \dots(6.39)$$

Multiplying this Equation by  $\left( \frac{\alpha}{2^n n! \sqrt{\pi}} \right)^{1/2} e^{-\alpha^2 x^2 / 2}$

and simplifying, we obtain

$$x \psi_n(x) = \frac{1}{\alpha \sqrt{2}} [\sqrt{n+1} \psi_{n+1}(x) + \sqrt{n} \psi_{n-1}(x)] \quad \dots(6.40)$$

Substituting in (6.37),

$$\langle x \rangle = \frac{\sqrt{n+1}}{\alpha \sqrt{2}} \int_{-\infty}^{\infty} \psi_n^*(x) \psi_{n+1}(x) dx + \frac{\sqrt{n}}{\alpha \sqrt{2}} \int_{-\infty}^{\infty} \psi_n^*(x) \psi_{n-1}(x) dx$$

Since the oscillator wave functions are orthonormal, both the integrals on the right-hand side vanish. Therefore,

$$\langle x \rangle = 0 \quad \dots(6.41)$$

$$\begin{aligned} \text{(b) } \langle x^2 \rangle &= \int_{-\infty}^{\infty} \psi_n^*(x) x^2 \psi_n(x) dx \\ &= \frac{1}{\sqrt{2} \alpha} \int_{-\infty}^{\infty} \psi_n^* x [\sqrt{n+1} \psi_{n+1} + \sqrt{n} \psi_{n-1}] dx \quad \text{Using (6.40)} \\ &= \frac{1}{\sqrt{2} \alpha} \left[ \sqrt{n+1} \int_{-\infty}^{\infty} \psi_n^* x \psi_{n+1} dx + \sqrt{n} \int_{-\infty}^{\infty} \psi_n^* x \psi_{n-1} dx \right] \end{aligned}$$

Using (6.40) again,

$$\begin{aligned}
\langle x^2 \rangle &= \frac{1}{\alpha} \sqrt{\frac{n+1}{2}} \left[ \int_{-\infty}^{\infty} \psi_n^* \left\{ \frac{1}{\alpha} \sqrt{\frac{n+2}{2}} \psi_{n+2} + \frac{1}{\alpha} \sqrt{\frac{n+1}{2}} \psi_n \right\} dx \right] \\
&+ \frac{1}{\alpha} \sqrt{\frac{n}{2}} \left[ \int_{-\infty}^{\infty} \psi_n^* \left\{ \frac{1}{\alpha} \sqrt{\frac{n}{2}} \psi_n + \frac{1}{\alpha} \sqrt{\frac{n-1}{2}} \psi_{n-2} \right\} dx \right] \\
&= \frac{\sqrt{(n+1)(n+2)}}{2\alpha^2} \int_{-\infty}^{\infty} \psi_n^* \psi_{n+2} dx + \frac{n+1}{2\alpha^2} \int_{-\infty}^{\infty} \psi_n^* \psi_n dx \\
&+ \frac{n}{2\alpha^2} \int_{-\infty}^{\infty} \psi_n^* \psi_n dx + \frac{\sqrt{n(n-1)}}{2\alpha^2} \int_{-\infty}^{\infty} \psi_n^* \psi_{n-2} dx
\end{aligned}$$

Using the orthonormality of oscillator wave functions, we obtain

$$\langle x^2 \rangle = \frac{1}{2\alpha^2} (n+1+n)$$

or 
$$\boxed{\langle x^2 \rangle = \frac{2n+1}{2\alpha^2} = \left(n + \frac{1}{2}\right) \frac{\hbar}{m\omega}} \quad \dots(6.42)$$

$$\begin{aligned}
\text{(c) } \langle p \rangle &= \int_{-\infty}^{\infty} \psi_n^*(x) \hat{p} \psi_n(x) dx \\
&= \int_{-\infty}^{\infty} \psi_n^*(x) \left[ -i\hbar \frac{d\psi_n(x)}{dx} \right] dx \quad \dots(6.43)
\end{aligned}$$

Now, if  $\psi_n(x)$  is odd, then its derivative is even, and vice versa. Therefore, the integrand in the above integral is always an odd function of  $x$ . Hence  $\langle p \rangle = 0$ .

We can also obtain this result using recurrence relation (6.23a):

$$\frac{dH_n(\alpha x)}{d(\alpha x)} = 2n H_{n-1}(\alpha x) \quad \dots(6.44)$$

Differentiating (6.39),

$$\begin{aligned}
\frac{d\psi_n(x)}{dx} &= \left( \frac{\alpha}{2^n n! \sqrt{\pi}} \right)^{1/2} \times \left[ -(\alpha^2 x) \exp\left( \frac{-\alpha^2 x^2}{2} \right) H_n(\alpha x) + \exp\left( \frac{-\alpha^2 x^2}{2} \right) \frac{dH_n(\alpha x)}{d(\alpha x)} \right] \\
&= -\alpha^2 x \left( \frac{\alpha}{2^n n! \sqrt{\pi}} \right)^{1/2} \exp\left( \frac{-\alpha^2 x^2}{2} \right) H_n(\alpha x) + \left( \frac{\alpha}{2^n n! \sqrt{\pi}} \right)^{1/2} \times \alpha \exp\left( \frac{-\alpha^2 x^2}{2} \right) \frac{dH_n(\alpha x)}{d(\alpha x)}
\end{aligned}$$

Using (6.39) and (6.44) this becomes

$$\begin{aligned}
\frac{d\psi_n(x)}{dx} &= -\alpha^2 x \psi_n(x) + \left( \frac{\alpha}{2^n n! \sqrt{\pi}} \right)^{1/2} \alpha \exp\left( \frac{-\alpha^2 x^2}{2} \right) [2n H_{n-1}(\alpha x)] \\
&= -\alpha^2 x \psi_n(x) + (2\alpha n) \frac{1}{\sqrt{2n}} \left( \frac{\alpha}{2^{n-1} (n-1)! \sqrt{\pi}} \right)^{1/2} \exp\left( \frac{-\alpha^2 x^2}{2} \right) H_{n-1}(\alpha x) \\
&= -\alpha^2 x \psi_n(x) + 2\alpha \left( \frac{n}{2} \right)^{1/2} \psi_{n-1}(x)
\end{aligned}$$

Using (6.40),

$$\frac{d\psi_n(x)}{dx} = \frac{-\alpha}{\sqrt{2}} [\sqrt{n+1} \psi_{n+1}(x) + \sqrt{n} \psi_{n-1}(x)] + \frac{2\alpha}{\sqrt{2}} \sqrt{n} \psi_{n-1}(x)$$

Or,

$$\frac{d\psi_n(x)}{dx} = \frac{\alpha}{\sqrt{2}} [\sqrt{n} \psi_{n-1}(x) - \sqrt{n+1} \psi_{n+1}(x)] \quad \dots(6.45)$$

Substituting in (6.43) and using the orthonormality of eigenfunctions, we get

$$\langle p \rangle = 0 \quad \dots(6.46)$$

$$\begin{aligned}
\text{(d) } \langle p^2 \rangle &= \int_{-\infty}^{\infty} \psi_n^*(x) \hat{p}^2 \psi_n(x) dx \\
&= -\hbar^2 \int_{-\infty}^{\infty} \psi_n^*(x) \frac{d^2 \psi_n(x)}{dx^2} dx \\
&= \frac{-\hbar^2 \alpha}{\sqrt{2}} \int_{-\infty}^{\infty} \psi_n^* \frac{d}{dx} [\sqrt{n} \psi_{n-1} - \sqrt{n+1} \psi_{n+1}] dx \quad \text{Using (6.45)} \\
&= \frac{-\hbar^2 \alpha}{\sqrt{2}} \left[ \sqrt{n} \int_{-\infty}^{\infty} \psi_n^* \frac{d\psi_{n-1}}{dx} dx - \sqrt{n+1} \int_{-\infty}^{\infty} \psi_n^* \frac{d\psi_{n+1}}{dx} dx \right]
\end{aligned}$$

Using (6.45) again,

$$\begin{aligned}
\langle p^2 \rangle &= \frac{-\hbar^2 \alpha^2 \sqrt{n}}{2} \int_{-\infty}^{\infty} \psi_n^* [\sqrt{n-1} \psi_{n-2}(x) - \sqrt{n} \psi_n] dx \\
&\quad + \frac{\hbar^2 \alpha^2 \sqrt{n+1}}{2} \int_{-\infty}^{\infty} \psi_n^* [\sqrt{n+1} \psi_n - \sqrt{n+2} \psi_{n+2}] dx \\
&= \frac{-\hbar^2 \alpha^2}{2} \left[ \sqrt{n(n-1)} \int_{-\infty}^{\infty} \psi_n^* \psi_{n-2} dx - n \int_{-\infty}^{\infty} \psi_n^* \psi_n dx \right. \\
&\quad \left. - (n+1) \int_{-\infty}^{\infty} \psi_n^* \psi_n dx + \sqrt{(n+1)(n+2)} \int_{-\infty}^{\infty} \psi_n^* \psi_{n+2} dx \right]
\end{aligned}$$

Using the orthonormality of oscillator wave function, this reduces to

$$\begin{aligned}
 \langle p^2 \rangle &= -\frac{\hbar^2 \alpha^2}{2} (0 - n - (n + 1) + 0) \\
 &= \frac{2n + 1}{2} \hbar^2 \alpha^2 \\
 &= \left( n + \frac{1}{2} \right) \hbar^2 \left( \frac{m\omega}{\hbar} \right) \\
 \boxed{\langle p^2 \rangle} &= \boxed{\left( n + \frac{1}{2} \right) m\omega\hbar} \quad \dots(6.47)
 \end{aligned}$$

**Problem 2:** Evaluate the position-momentum uncertainty product for the  $n$ th state of a linear harmonic oscillator.

**Solution:**

$$\begin{aligned}
 (\Delta x)^2 &= \langle x^2 \rangle - \langle x \rangle^2 \\
 &= \left( n + \frac{1}{2} \right) \frac{\hbar}{m\omega} - 0 \\
 &= \left( n + \frac{1}{2} \right) \frac{\hbar}{m\omega} \\
 (\Delta p)^2 &= \langle p^2 \rangle - \langle p \rangle^2 \\
 &= \left( n + \frac{1}{2} \right) m\omega\hbar - 0 = \left( n + \frac{1}{2} \right) m\omega\hbar
 \end{aligned}$$

Multiplying the two and taking square root we obtain the position-momentum uncertainty product

$$\Delta x \Delta p = \left( n + \frac{1}{2} \right) \hbar, \quad n = 0, 1, 2, \dots \quad \dots(6.48)$$

This is in accordance with the uncertainty relation

$$\Delta x \Delta p \geq \frac{\hbar}{2}$$

For  $n = 0$ , (6.48) reduces to

$$\Delta x \Delta p = \frac{\hbar}{2}$$

showing that *the uncertainty product is minimum for the ground state.*

**Problem 3:** Obtain the expectation values of the kinetic and potential energies for the  $n$ th state of a linear harmonic oscillator.

$$\text{Solution: } \langle T \rangle = \frac{1}{2m} \langle p^2 \rangle = \frac{1}{2} \left( n + \frac{1}{2} \right) \hbar\omega = \frac{E_n}{2}$$

$$\begin{aligned}\langle V \rangle &= \frac{1}{2} k \langle x^2 \rangle = \frac{1}{2} (m\omega^2) \left( n + \frac{1}{2} \right) \frac{\hbar}{m\omega} \\ &= \frac{1}{2} \left( n + \frac{1}{2} \right) \hbar\omega = \frac{E_n}{2}\end{aligned}$$

Thus, the average kinetic and potential energies for a harmonic oscillator in any eigenstate are each equal to one-half the total energy, as in the case of a classical harmonic oscillator.

**H.W.:** Write down the required information for the linear harmonic oscillator in the state  $\psi_2(x)$ ;

- 1- Figure for  $\psi_2(x)$ ; as a function of  $x$ .
- 2- The parity of  $\psi_2(x)$ ;
- 3- The number of nodes.
- 4- The energy state  $E_2$ .

**H.W.:** Find the expectation value of the kinetic energy for the Linear harmonic oscillator in the eigen energy state  $n = 0$ .

**H.W.:** Prove that the following wave functions are orthonormal functions.

$$\begin{aligned}\psi_0(x) &= \left( \frac{\alpha}{\sqrt{\pi}} \right)^{1/2} \exp\left(-\frac{1}{2} \alpha^2 x^2\right) \\ \psi_1(x) &= \left( \frac{2\alpha}{\sqrt{\pi}} \right)^{1/2} (\alpha x) \exp\left(-\frac{1}{2} \alpha^2 x^2\right) \\ \psi_2(x) &= \left( \frac{\alpha}{2\sqrt{\pi}} \right)^{1/2} (2\alpha^2 x^2 - 1) \exp\left(-\frac{1}{2} \alpha^2 x^2\right)\end{aligned}$$

**H.W.:** Use the uncertainty principle to prove that the minimum energy for the linear harmonic oscillator is  $\frac{1}{2} \hbar\omega$ .