

3.7 Probability Conservation and the Hermiticity of the Hamiltonian

If the operator \hat{A} satisfy the condition
is called Hermitian operator.

$$\int \Psi^* \hat{A} \Psi d\mathbf{r} = \int (\hat{A} \Psi)^* \Psi d\mathbf{r}$$

We shall now show that the conservation of probability implies that the Hamiltonian operator \mathbf{H} appearing in the Schrödinger equation is **Hermitian**.

In terms of \mathbf{H} , the Schrödinger equation can be written as

$$i\hbar \frac{\partial \Psi}{\partial t} = H\Psi \quad \dots(3.30)$$

The complex conjugate of this equation is

$$-i\hbar \frac{\partial \Psi^*}{\partial t} = (H\Psi)^* \quad \dots(3.31)$$

Using these equations, we can write

$$\begin{aligned} \frac{\partial}{\partial t} \int \Psi^* \Psi d\mathbf{r} &= \int \left(\Psi^* \frac{\partial \Psi}{\partial t} + \Psi \frac{\partial \Psi^*}{\partial t} \right) d\mathbf{r} \quad \dots(3.32) \\ &= (i\hbar)^{-1} \int [\Psi^* (H\Psi) - \Psi (H\Psi)^*] d\mathbf{r} \end{aligned}$$

Since the left-hand side is zero, we obtain

$$\int \Psi^* (H\Psi) d\mathbf{r} = \int (H\Psi)^* \Psi d\mathbf{r} \quad \dots(3.33)$$

Operators which satisfy this condition are called Hermitian. Thus, H is an Hermitian operator.

3.8 Probability Current Density

From equation (3.28)

$$\frac{\partial}{\partial t} \int_V P(\mathbf{r}, t) d\mathbf{r} = - \int_V \nabla \cdot \mathbf{j} d\mathbf{r}$$

and since this equation is true for any arbitrary volume, we have

$$\frac{\partial}{\partial t} P(\mathbf{r}, t) + \nabla \cdot \mathbf{j}(\mathbf{r}, t) = 0 \quad \dots(3.34)$$

This is called the **equation of continuity**.

The vector $\mathbf{j}(\mathbf{r}, t)$ is called the **probability current density**.

If $\nabla \cdot \mathbf{j} = 0$, then for that state the probability density is constant in time. Such states are called **stationary states**.

3.8 Expectation Values of Dynamical Variables

- In quantum mechanics a particle is represented by a wave function which can be obtained by solving the Schrödinger equation and contains all the available information about the particle.
- The **dynamical variables** of the particle (position, momentum, ...) can be extracted from the wave function Ψ .
- Since Ψ has a probabilistic interpretation, so that exact information about the variables cannot be obtained.
- We obtain only the **expectation value** of a quantity, which is **the average of a large number of measurements on the same system**.
- The expectation value of a physical quantity is always real.

To find the position expectation value $\langle r \rangle$:

Since $P(\mathbf{r}, t) = \Psi^*(\mathbf{r}, t) \Psi(\mathbf{r}, t)$ is interpreted as the position probability density at the point \mathbf{r} at the time t , the *expectation value* of the position vector \mathbf{r} is given by

$$\langle \mathbf{r} \rangle = \int \mathbf{r} P(\mathbf{r}, t) d\mathbf{r} \quad \dots(3.35)$$

$$= \int \Psi^*(\mathbf{r}, t) \mathbf{r} \Psi(\mathbf{r}, t) d\mathbf{r} \quad \dots(3.36)$$

where $\Psi(\mathbf{r}, t)$ is normalized. This equation is equivalent to the three equations

]

$$\begin{aligned}\langle x \rangle &= \int \Psi^* x \Psi \, d\mathbf{r} \\ \langle y \rangle &= \int \Psi^* y \Psi \, d\mathbf{r} \\ \langle z \rangle &= \int \Psi^* z \Psi \, d\mathbf{r}\end{aligned}\quad \dots(3.37)$$

- The expectation value is a function only of the time because the space coordinates have been integrated out.
- The expectation value of a physical quantity is always real.

The expectation value of any quantity which is a function of \mathbf{r} and t would be;

$$\boxed{\langle f(\mathbf{r}, t) \rangle = \int \Psi^*(\mathbf{r}, t) f(\mathbf{r}, t) \Psi(\mathbf{r}, t) \, d\mathbf{r}} \quad \dots(3.38)$$

As an example, the expectation value of the potential energy is

$$\langle V(\mathbf{r}, t) \rangle = \int \Psi^*(\mathbf{r}, t) V(\mathbf{r}, t) \Psi(\mathbf{r}, t) \, d\mathbf{r} \quad \dots(3.39)$$

The expectation values for quantities which are functions of momentum or of both position and momentum by using the *operator representations*

$$\begin{aligned}\hat{\mathbf{p}} &= -i\hbar \nabla \\ p^2 &= -\hbar^2 \nabla^2 \\ \hat{E} &= i\hbar \frac{\partial}{\partial t}\end{aligned}$$

Q: How these differential operators are to be combined with the position probability density $\Psi^* \Psi$ to obtain the total energy expressions?

Answer: by using the classical expression for the total energy

$$E = \frac{p^2}{2m} + V$$

And the expectation values;

$$\langle E \rangle = \left\langle \frac{p^2}{2m} \right\rangle + \langle V \rangle \quad \dots(3.40)$$

Replacing E and p^2 by the corresponding operators, we get

$$\left\langle i\hbar \frac{\partial}{\partial t} \right\rangle = \left\langle -\frac{\hbar^2}{2m} \nabla^2 \right\rangle + \langle V \rangle \quad \dots(3.41)$$

This equation must be consistent with the Schrödinger equation

$$i\hbar \frac{\partial \Psi}{\partial t} = -\frac{\hbar^2 \nabla^2}{2m} \Psi + V\Psi$$

Multiplying by Ψ^* on the left and integrating, we get

$$\int \Psi^* \left(i\hbar \frac{\partial}{\partial t} \right) \Psi d\mathbf{r} = \int \Psi^* \left(-\frac{\hbar^2 \nabla^2}{2m} \right) \Psi d\mathbf{r} + \int \Psi^* V \Psi d\mathbf{r} \quad \dots(3.42)$$

From (3.40) -(3.42) we have;

$$\langle V \rangle = \int \Psi^* V \Psi d\mathbf{r} \quad \dots(3.43)$$

$$\langle E \rangle = \int \Psi^* i\hbar \frac{\partial \Psi}{\partial t} d\mathbf{r} \quad \dots(3.44)$$

$$\langle \mathbf{p} \rangle = \int \Psi^* (-i\hbar) \nabla \Psi d\mathbf{r} \quad \dots(3.45)$$

The last equation is equivalent to

$$\begin{aligned} \langle p_x \rangle &= -i\hbar \int \Psi^* \frac{\partial \Psi}{\partial x} d\mathbf{r} \\ \langle p_y \rangle &= -i\hbar \int \Psi^* \frac{\partial \Psi}{\partial y} d\mathbf{r} \\ \langle p_z \rangle &= -i\hbar \int \Psi^* \frac{\partial \Psi}{\partial z} d\mathbf{r} \end{aligned} \quad \left. \vphantom{\begin{aligned} \langle p_x \rangle \\ \langle p_y \rangle \\ \langle p_z \rangle \end{aligned}} \right] \quad \dots(3.46)$$

Generalizing the above results:

$\Psi(\mathbf{r}, t)$ The normalized wave function which describes the dynamical state of a particle.

$A(\mathbf{r}, \mathbf{p}, t)$ the dynamical variable representing a physical quantity associated with the particle.

We calculate the expectation value of A from the expression;

$$\langle A \rangle = \int \Psi^*(\mathbf{r}, t) \hat{A} \Psi(\mathbf{r}, t) d\mathbf{r}, \quad \text{or}$$

$$\boxed{\langle A \rangle = \int \Psi^*(\mathbf{r}, t) \hat{A}(\mathbf{r}, -i\hbar\nabla, t) \Psi(\mathbf{r}, t) d\mathbf{r}} \quad \dots(3.46)$$

Since the *expectation value of a physical quantity is always real*, i.e., $\langle A \rangle^* = \langle A \rangle$

Then the operator \hat{A} must satisfy;

$$\boxed{\int \Psi^* \hat{A} \Psi d\mathbf{r} = \int (\hat{A}\Psi)^* \Psi d\mathbf{r}} \quad \dots(3.47)$$

Thus, the operator associated with a dynamical quantity must be **Hermitian**.