## Chapter Three

## Steady state Heat Conduction in Plane Walls

1) Considerable temperature difference between the inner and the outer surfaces of the wall (significant temperature gradient in the $x$ direction).
2) The wall surface is nearly isothermal.


Assuming heat transfer is the only energy interaction and there is no heat generation, the energy balance can be expressed as

Rate of change of the energy of the wall

$$
\dot{Q}_{\text {in }}-\dot{Q}_{\text {out }}=\frac{d E_{\text {wall }}}{d t}=0
$$

Then Fourier's law of heat conduction for the wall can be expressed as

$$
\begin{equation*}
\dot{Q}_{\text {cond }, \text { wall }}=-k A \frac{d T}{d x} \tag{W}
\end{equation*}
$$

Integrating the above equation and rearranging yields

$$
\dot{Q}_{\text {cond,wall }}=k A \frac{T_{1}-T_{2}}{L} \quad(\mathrm{~W})
$$

## Conduction Resistance

The above equation for heat conduction through a plane wall can be rearranged as

$$
\begin{equation*}
\dot{Q}_{\text {cond }, \text { wall }}=\frac{T_{1}-T_{2}}{R_{\text {wall }}} \tag{W}
\end{equation*}
$$

Where $R_{\text {wall }}$ is the conduction resistance expressed as

$$
R_{\text {wall }}=\frac{L}{k A} \quad\left({ }^{\circ} \mathrm{C} / \mathrm{W}\right)
$$

## Convection Resistance

Newton's law of cooling for convection heat transfer rate $\left(\dot{Q}_{\text {conv }}=h A_{s}\left(T_{s}-T_{\infty}\right)\right)$ can be rearranged as

$$
\begin{equation*}
\dot{Q}_{c o n v}=\frac{T_{s}-T_{\infty}}{R_{c o n v}} \tag{W}
\end{equation*}
$$

$R_{\text {conv }}$ is the convection resistance

$$
R_{c o n v}=\frac{1}{h A_{s}} \quad\left({ }^{\circ} \mathrm{C} / \mathrm{W}\right)
$$



## Radiation Resistance

$$
\begin{aligned}
\dot{Q}_{r a d} & =\varepsilon \sigma A_{s}\left(T_{s}^{4}-T_{\text {surr }}^{4}\right)=h_{r a d} A_{s}\left(T_{s}-T_{\text {surr }}\right)=\frac{T_{s}-T_{\text {surr }}}{R_{\text {rad }}}(\mathrm{W}) \\
R_{r a d} & =\frac{1}{h_{r a d} A_{s}} \quad(K / \mathrm{W}) \\
h_{r a d} & =\frac{\dot{Q}_{\text {rad }}}{A_{s}\left(T_{s}-T_{\text {surr }}\right)}=\varepsilon \sigma\left(T_{s}^{2}+T_{\text {surr }}^{2}\right)\left(T_{s}+T_{\text {surr }}\right)\left(\mathrm{W} / \mathrm{m}^{2} \cdot \mathrm{~K}\right)
\end{aligned}
$$

## Radiation and Convection Resistance

$$
h_{\text {combined }}=h_{\text {conv }}+h_{\text {rad }}
$$

## Thermal Resistance Network

$$
\begin{aligned}
& \dot{Q}=h_{1} A\left(T_{\infty, 1}-T_{1}\right)= \\
& k A \frac{T_{1}-T_{2}}{L}=h_{2} A\left(T_{2}-T_{\infty, 2}\right)
\end{aligned}
$$



$$
R_{\text {total }}=R_{\text {conv }, 1}+R_{\text {wall }}+R_{\text {conv }, 2}=\frac{1}{h_{1} A}+\frac{L}{k A}+\frac{1}{h_{2} A}\left({ }^{\circ} \mathrm{C} / \mathrm{W}\right)
$$

## Multilayer Plane Walls

$$
\begin{aligned}
R_{\text {total }} & =R_{\text {conv }, 1}+R_{\text {wall }, 1}+R_{\text {wall }, 2}+R_{\text {conv }, 2} \\
& =\frac{1}{h_{1} A}+\frac{L_{1}}{k_{1} A}+\frac{L_{2}}{k_{2} A}+\frac{1}{h_{2} A}
\end{aligned}
$$

## Generalized Thermal Resistance Networks

The thermal resistance network, which consists of two parallel resistances, can be represented as shown in the figure. Noting that the total heat transfer is the sum of the heat transfers through each layer, we have

$$
\dot{Q}=\dot{Q}_{1}+\dot{Q}_{2}=\frac{T_{1}-T_{2}}{R_{1}}+\frac{T_{1}-T_{2}}{R_{2}}=\left(T_{1}-T_{2}\right)\left(\frac{1}{R_{1}}+\frac{1}{R_{2}}\right)
$$

Utilizing electrical analogy, we get
$\dot{Q}=\frac{T_{1}-T_{2}}{R_{\text {total }}} \quad$ where $\quad \frac{1}{R_{\text {total }}}=\frac{1}{R_{1}}+\frac{1}{R_{2}} \rightarrow R_{\text {total }}=\frac{R_{1} R_{2}}{R_{1}+R_{2}}$


Now consider the combined series-parallel arrangement shown, the total rate of heat transfer through this composite system can again be expressed as,

$$
\text { where } \dot{Q}=\frac{\bar{T}_{1}-T_{\infty}}{R_{\text {total }}}
$$

$$
\text { and } \quad R_{\text {total }}=R_{12}+R_{3}+R_{\mathrm{conv}}=\frac{R_{1} R_{2}}{R_{1}+R_{2}}+R_{3}+R_{\mathrm{conv}}
$$

$$
R_{1}=\frac{L_{1}}{k_{1} A_{1}}, \quad R_{2}=\frac{L_{2}}{k_{2} A_{2}}, \quad R_{3}=\frac{L_{3}}{k_{3} A_{3}}, \quad R_{\mathrm{conv}}=\frac{1}{h A_{3}}
$$

once the individual thermal resistances are evaluated, the total resistance and the total rate of heat transfer can easily be determined from the relations above.


## Sphere

Sphere systems may also be treated as on dimensional when the temperature is a function of radius only

$$
q=\frac{4 \pi k\left(T_{i}-T_{o}\right)}{1 / r_{i}-1 / r_{o}}
$$

## Heat Conduction in Cylinders

$$
\begin{align*}
& \dot{Q}_{\text {cond }, c y l}=-k A \frac{d T}{d r} \quad(\mathrm{~W})  \tag{W}\\
& \int_{r=r_{1}}^{r_{2}} \frac{\dot{Q}_{\text {cond }, c y l}}{A} d r=-\int_{T=T_{1}}^{T_{2}} k d T \\
& \dot{Q}_{\text {cond }, c y l}=2 \pi L k \frac{T_{1}-T_{2}}{\ln \left(r_{2} / r_{1}\right)}
\end{align*}
$$



Thermal Resistance with Convection

$$
\begin{aligned}
& \dot{Q}_{\text {cond }, c y l}=\frac{T_{1}-T_{2}}{R_{c y l}} \\
& \dot{Q}=\frac{T_{\infty, 1}-T_{\infty, 2}}{R_{\text {total }}} \\
& R_{\text {total }}=R_{\text {conv }, 1}+R_{\text {cyl }}+R_{\text {conv }, 2}= \\
& \quad=\frac{1}{\left(2 \pi r_{1} L\right) h_{1}}+\frac{\ln \left(r_{2} / r_{1}\right)}{2 \pi L k}+\frac{1}{\left(2 \pi r_{2} L\right) h_{2}}
\end{aligned}
$$




## Multilayered Cylinders

$$
\begin{aligned}
& R_{\text {total }}=R_{c o n v, 1}+R_{c y l, 1}+R_{c y l, 3}+R_{c y l, 3}+R_{c o n v, 2}= \\
& =\frac{1}{\left(2 \pi r_{1} L\right) h_{1}}+\frac{\ln \left(r_{2} / r_{1}\right)}{2 \pi L k_{1}}+\frac{\ln \left(r_{3} / r_{2}\right)}{2 \pi L k_{2}}+\frac{\ln \left(r_{4} / r_{3}\right)}{2 \pi L k_{3}}+\frac{1}{\left(2 \pi r_{2} L\right) h_{2}}
\end{aligned}
$$



