

# Structural, optical and electrical properties of $\text{Cu}_2\text{Zn}_{1-x}\text{Cd}_x\text{SnS}_4$ quinary alloys nanostructures deposited on porous silicon

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**Abstract** The  $\text{Cu}_2\text{Zn}_{1-x}\text{Cd}_x\text{SnS}_4$  quinary alloy nanostructures with different Cd contents were grown using spin coating technique on porous silicon (63.93 %) substrate. The structural properties of  $\text{Cu}_2\text{Zn}_{1-x}\text{Cd}_x\text{SnS}_4/\text{PS}$  were investigated by X-ray diffraction and field emission-scanning electron microscope (FE-SEM). The optical properties studied through photoluminescence technique, indicated that the band gap is shifted as Cd content increases from 1.84 eV at  $x = 0$  to 1.76 eV at  $x = 1$ . The electrical characterization of the Ag/n-PS/ $\text{Cu}_2\text{Zn}_{1-x}\text{Cd}_x\text{SnS}_4/\text{Ag}$  diode through current to voltage (I–V) characterization shows the highest photo-response of (value if any) at  $\text{Cu}_2\text{Zn}_{0.4}\text{Cd}_{0.6}\text{SnS}_4$  composition.

## 1 Introduction

A kesterite structure type of  $\text{Cu}_2\text{ZnSnS}_4$  (CZTS) is a potential absorber material as thin film solar cells due to presence of naturally abundant, inexpensive, efficient and environment friendly elements. It is known that CZTS has a direct band gap between 1.45 and 1.6 eV with high absorption coefficient over  $10^4 \text{ cm}^{-1}$  (in visible light), intrinsic p-type conductivity and low thermal conductivity (Katagiri et al. 2001; Kamoun et al. 2007; Levenco et al. 2012). These properties promote CZTS as a good candidate for photovoltaic materials. The predicted theoretical efficiency of CZTS based solar cell is larger than 30 % (Daranfed et al. 2012; Todorov et al. 2010; Guo et al. 2009; Katagiri et al. 2009). Efficiency up to 6.77 % has been already reached in solar cells produced with CZTS as absorber layer (Canham 1997). Silicon substrates have high quality single crystal with nearly defect free, very low volume density of impurities and controlled amount of dopants. It may be prepared flat on the atomic scale by standard methods that lead to porous silicon (PS) formation (Yerokhov and Melnyk 1999). The PS is attractive for solar cell applications due to its efficient antireflection coatings and other properties such as band gap broadening (Schirone et al. 1997), wide absorption spectrum (Menna et al. 1997) and optical transmission range (700–1000 nm) (Schirone et al. 1998). It can also be used for surface passivation and texturization (Cláudia et al. 2008).

Al-Douri et al. (2011) have calculated the energy gap of Si that is found to be indirect using empirical pseudo potential method (EPM). They have investigated features such as refractive index, optical dielectric constant, bulk modulus, elastic constants and short-range force constants, in addition to the shear modulus, Young's modulus, Poisson's ratio

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and Lamé's constants for both Si and PS. The Debye temperature of PS was estimated from the average sound velocity.  $\text{Cu}_2\text{ZnSnS}_4$  (CZTS) film was fabricated by sulfurization of two stacked layers of Cu(Zn,Sn) (CZT) alloy precursors by Hong and Kim (2013). The sulfurization was performed in an evacuated and sealed quartz ampoule with sulphur powder, in which samples were annealed at 450, 500, or 550 °C for 1 h. The energy band gap of CZTS absorber was determined to be 1.370–1.414 eV. While, Xinkun et al. (2012) have prepared Sn/Cu/ZnS by evaporation on soda lime glass at room temperature, and then polycrystalline thin films of  $\text{Cu}_2\text{ZnSnS}_4$  (CZTS) were produced by sulfurizing the precursors in a sulfur atmosphere at a temperature of 550 °C for 3 h. The experimental results show that, when the ratios of  $[\text{Cu}]/([\text{Zn}] + [\text{Sn}])$  and  $[\text{Zn}]/[\text{Sn}]$  in the CZTS are 0.83 and 1.15, the CZTS thin films possess an absorption coefficient larger than  $4.0 \times 10^4 \text{ cm}^{-1}$  in the energy range 1.5–3.5 eV, and a direct band gap (1.47 eV). Therefore, the CZTS thin films are suitable for absorption layers of solar cells. Accordingly, successful experiments have discovered that electrochemical and chemical dissolution enable the silicon wafers to emit light in red luminescence (Canham 1990). During the last years, the optical properties of PS have become a very intense area of research.

The prime aim of this work is to study the structural properties given by X-ray diffraction (XRD) and field emission-scanning electron microscope (FE-SEM), and optical properties of  $\text{Cu}_2\text{Zn}_{1-x}\text{Cd}_x\text{SnS}_4$  quaternary alloy nanostructures by PL, in addition to investigate the electrical properties of Ag/n-PS/ $\text{Cu}_2\text{Zn}_{1-x}\text{Cd}_x\text{SnS}_4$ /Ag diode at  $x = 0, 0.6, 1$  for photodetector applications through (I–V) characterization. The paper is organized as the followings; Sect. 2 details the experimental procedure. The results and discussion are elaborated in Sect. 3. Finally, conclusion is summarised in Sect. 4.

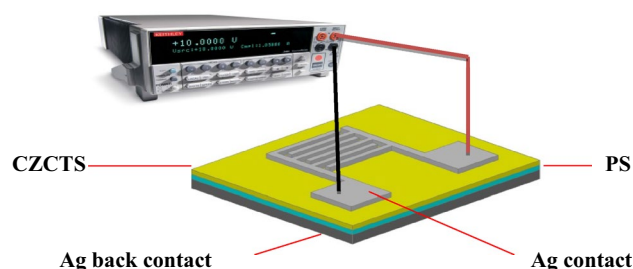
## 2 Experimental

**Spin coating** technique is used to deposit the  $\text{Cu}_2\text{Zn}_{1-x}\text{Cd}_x\text{SnS}_4$  quaternary alloy nanostructures onto **n-type** PS substrate. The n-PS substrate was prepared as reported by Abd et al. (2013) where the optimal case (63.93 %) was selected. First the solution of  $\text{Cu}_2\text{Zn}_{1-x}\text{Cd}_x\text{SnS}_4$  precursors was prepared by **dissolving copper (II) chloride dihydrate (0.3 M), zinc (II) chloride dihydrate (0.3 M), tin (II) chloride dehydrate (0.3 M), cadmium (II) chloride (0.3 M), thiourea (0.6)** was dissolved in 2-methoxyethanol and **monoethanolamine (MEA—as a stabilizer)** was added as a stabilizer. The mole ratios of Cu, (Zn + Cd), Sn, and S in the solution were **2:1:1:4**. For obtaining the solution with different Cd content ( $x$ ), the mole ratios of Cd to (Zn + Cd) in the solution vary according to the value of  $x$  as 0, 0.6,

and 1. The solution was **stirred at 50 °C** to completely dissolve the metal compounds during stirring **the milk solution became yellow in colour**. Afterthat, the solution was dropped onto optimal etching substrate, which was rotating at 2500 rpm for 30 s. After deposition by spin coating, the nanostructures were dried at 250 °C for 80 min on hot plate. The coating and drying processes were repeated for seven times to obtain 1  $\mu\text{m}$  thickness. The thickness measurement has been carried out using the weight method by:

$$t = \frac{\Delta m}{A \times \rho},$$

where  $t$  is thickness,  $\Delta m$  is difference of substrate weight (substrate after deposition – substrate before deposition),  $A$  is area of sample and  $\rho$  is density of deposited material. **Finally, heat treatment was conducted in an elevator furnace under  $\text{N}_2$  gas flow (5 % gas atmosphere) for 1 h at 300 °C and after annealing, the samples was cooled to below 40 °C in the chamber.** Afterthat, Ag metal contacts were formed on  $\text{Cu}_2\text{Zn}_{1-x}\text{Cd}_x\text{SnS}_4$  quaternary alloy nanostructures with Cd content equals 0, 0.6, 1. PVD-HANDY/2STE (Vaksis Company, USA) vacuum thermal evaporation in the pressure of  $4.5 \times 10^5$  Torr was used for deposition on oxidized silicon, and the contacts were formed in the form of zig-zag with 5 mm length and 100 nm thickness as shown in Fig. 1. The contact area of the diode was found to be  $3.14 \times 10^2 \text{ cm}^2$ . This work is analysed by XRD analysis as carried out on the as-deposited films to check the crystalline structure using (Philips PW 1710 X-ray diffractometer, USA), which record the intensity as a function of Bragg's angle in  $2\theta$ -range from 20–60° using Cu  $K\alpha$  ( $\lambda = 1.5406 \text{ \AA}$ ). Photoluminescence (PL) spectroscopy system (Perkin Elmer Lambda 950, USA) was using He-Cd laser with the wavelength ( $\lambda$ ) of 325 nm. Surface morphologies for  $\text{Cu}_2\text{Zn}_{1-x}\text{Cd}_x\text{SnS}_4$  quaternary alloy nanostructures were investigated by FE-SEM system (NOVA NANO SEM 450, USA). For the current to voltage (I–V) characterization, the fabricated device was connected in parallel with the Kiethly 2400 source meter, USA. The reading was recorded from –6 to 6 V. For current to time (I–t) analysis, the device was connected in series with the multimeter and

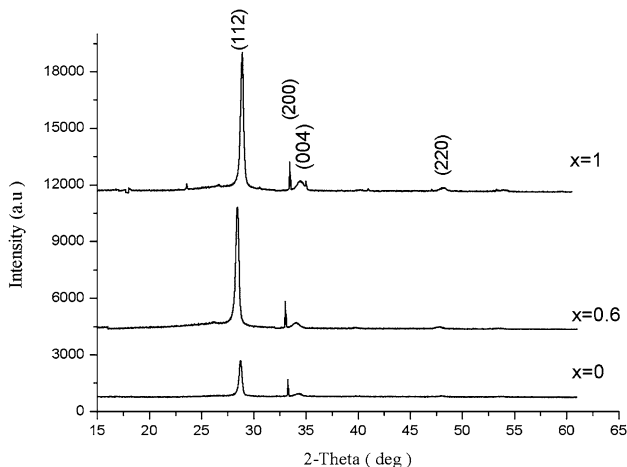


**Fig. 1** p–n junction

the value was recorded by switching the LED on and off. The wavelength and power of the LED used for conducting the experiment were 490 nm wavelength and 2 mW power, respectively.

### 3 Results and discussion

The XRD patterns of  $\text{Cu}_2\text{Cd}_x\text{Zn}_{1-x}\text{SnS}_4$  quinary alloy nanostructures with  $x = 0, 0.6$  and  $1$  deposited on PS prepared at the current density of  $5 \text{ mA/cm}^2$  for 120 min are shown in Fig. 2. The major diffraction peaks appear at  $2\theta = 28.757^\circ, 32.313^\circ, 34.275^\circ$  and  $47.30^\circ$  which can be attributed to the (112), (200), (004) and (220) planes of kesterite phase of  $\text{Cu}_2\text{ZnSnS}_4$  (ICDD PDF2008, 01-075-4122) and stannite phase of  $\text{Cu}_2\text{CdSnS}_4$  (ICDD PDF2008, 00-029-0537). It shows the highest (112) diffraction peak at  $x = 1$ . The peaks are shifted to the lower angle side with



**Fig. 2** XRD patterns of  $\text{Cu}_2\text{Zn}_{1-x}\text{Cd}_x\text{SnS}_4$  quinary alloy nanostructures at different Cd contents,  $x = 0, 0.6, 1$

increasing Cd content in the CZCTS solid solutions, which was attributed to the increasing lattice constant. This was because the radius of Cd ion ( $1.53 \text{ \AA}$ ) is larger than that of Zn ( $1.33 \text{ \AA}$ ) as supported by previous work (Ibraheem et al. 2015). The simplest possibility is that Cd substitutes other metals at their sites in crystal lattice of  $\text{Cu}_2\text{Cd}_x\text{Zn}_{1-x}\text{SnS}_4$  quinary alloy nanostructures. As the theoretical calculated substitution energies of Cd atoms at Cu, Sn and Zn atom sites in CZTS lattice are  $E_{\text{sub}}(\text{Cd}_{\text{Cu}}) = 0.69 \text{ eV}$ ,  $E_{\text{sub}}(\text{Cd}_{\text{Sn}}) = 1.07 \text{ eV}$  and  $E_{\text{sub}}(\text{Cd}_{\text{Zn}}) = 0.53 \text{ eV}$  (Maeda et al. 2012), the most likely is the isoelectronic substitution of Cd at the Zn site. Lattice constants  $a$  and  $c$  were calculated from XRD data for the (112) plane, which are given in Table 1;

$$\left(1/d^2\right) = \left(\left(h^2 + k^2\right) / a^2\right) + \left(l^2 / c^2\right), \quad (1)$$

where  $hkl$  are the Miller indices,  $a$  and  $c$  are the lattice constants. The interplane distance ( $d$ ) was calculated for all of CZCTS nanostructures using Bragg's diffraction equation (Ibraheem et al. 2015)

$$d = n\lambda / 2\sin\theta, \quad (2)$$

where  $\lambda$  is wavelength of XRD using ( $\lambda = 1.5406 \text{ \AA}$ ) and  $\theta$  is the Bragg's angle.

The crystallite size ( $D$ ) was calculated by Scherrer's formula (Ibraheem et al. 2015)

$$D = k\lambda / \beta\cos\theta. \quad (3)$$

where  $k$  is a constant, taken to be 0.94, and  $\beta$  is the full width at half maximum (FWHM).

It is known that the bulk modulus is a reflectance of the materials stiffness, which is important in different industries. Many authors (Sherry and Kumar 1991; Tallon 1980; Al-Douri et al. 2002, 2005; Al-Douri 2003, 2012) have made various efforts to explore thermodynamic properties of solids. In these studies, authors have examined the

**Table 1** The structural properties of  $\text{Cu}_2\text{Zn}_{1-x}\text{Cd}_x\text{SnS}_4$  quinary alloy nanostructures using XRD for different Cd contents,  $x = 0, 0.6, 1$

X	Peak ( $\theta$ )	Particle size (nm)	$D_{\text{hkl}}$ (112) ( $\text{\AA}$ )	Lattice constants ( $\text{\AA}$ )	Bulk modulus (GPa)
0	28.757	30.71	3.463	$a = 5.429$ $a = 5.427^a$ $c = 10.856$ $c = 10.848^a$	84.96
0.6	28.736	41.80	3.558	$a = 5.525$ $c = 10.906$	79.90
1	28.723	53.52	3.703	$a = 5.586$ $a = 5.487^b$ $c = 11.992$ $c = 11.389^c$	76.89

<sup>a</sup> Ref. (Ibraheem et al. 2015) exp.

<sup>b</sup> Ref. (Guan et al. 2013) exp.

<sup>c</sup> Ref. (Matsushita et al. 2000) exp.

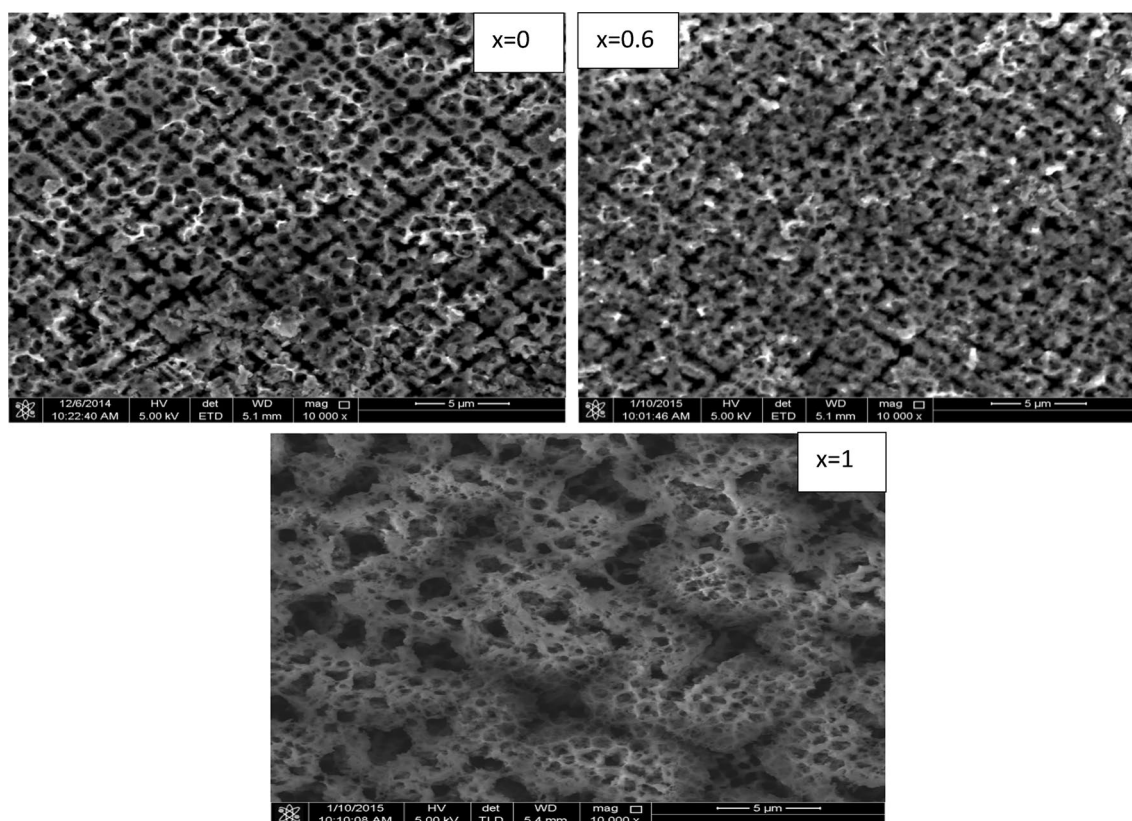
thermodynamic properties such as the inter-atomic separation and the bulk modulus of solids with different approximations and best-fit relations (Al-Douri et al. 2002, 2005; Al-Douri 2003, 2012). It has become possible to compute with great accuracy an important number of structural and electronic properties of solids. The ab initio calculations are complex and require significant effort. Therefore, more empirical approaches have been developed (Phillips 1973; Harison 1989) to compute properties of materials. In many cases, the empirical methods offer the advantage of applicability to a broad class of materials and to illustrate trends. In many applications, these empirical approaches do not give highly accurate results for each specific material, but are still very useful. Cohen (1985) has established an empirical formula for calculation of the bulk modulus  $B_0$ ; based on the nearest-neighbor distance. His result is in agreement with experimental values. Lam et al. (Lam et al. 1987) have derived an analytical expression for the bulk modulus from the total energy. This expression is different in structure from the empirical formula but gives similar numerical results. Also, they have obtained an analytical expression for the pressure derivative  $B_0$  of the bulk modulus. Our group (Al-Douri et al. 2004) used a concept based on the lattice constant to establish an empirical formula for the calculation of the bulk modulus. The results are in good

agreement with experimental data and other calculations. The theory yields a formula with two attractive features. Only the lattice constant is required as input, the computation of  $B_0$  itself is trivial. Consideration of hypothetical structure and simulation of the experimental conditions are required to make practical use of this formula.

The aim is to see how a qualitative concept, such as the bulk modulus, can be related to the lattice constant. It was argued that the dominant effect is the degree of covalency characterized by Phillips' homopolar gap  $E_h$  (Phillips 1973), and one reason for presenting these data in this work is that the validity of our calculations that is not restricted in computed space. We thus believe that the data will prove valuable for future work in this field. An important reason for studying  $B_0$  is the observation of clear differences between the lattice constants for CZCST quinary alloy nanostructures as seen in Table 1. The basis of our model is the lattice constant as seen in Table 1. Fitting of these data gives the following empirical formula (Al-Douri et al. 2004):

$$B_0 = [3000 - 100\lambda] \left(\frac{a}{2}\right)^{-3.5} \quad (4)$$

where  $a$  is the lattice constant (in Å) and  $\lambda$  is an empirical parameter which accounts for the effect of ionicity;  $\lambda = 0$ ;

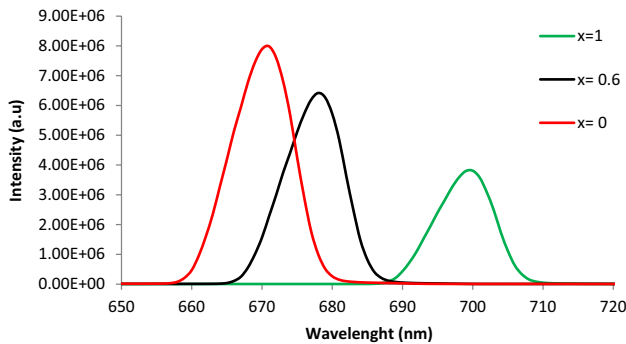


**Fig. 3** FE-SEM images of  $\text{Cu}_2\text{Zn}_{1-x}\text{Cd}_x\text{SnS}_4$  quinary alloy nanostructure grown on PS at different Cd contents,  $x = 0, 0.6, 1$

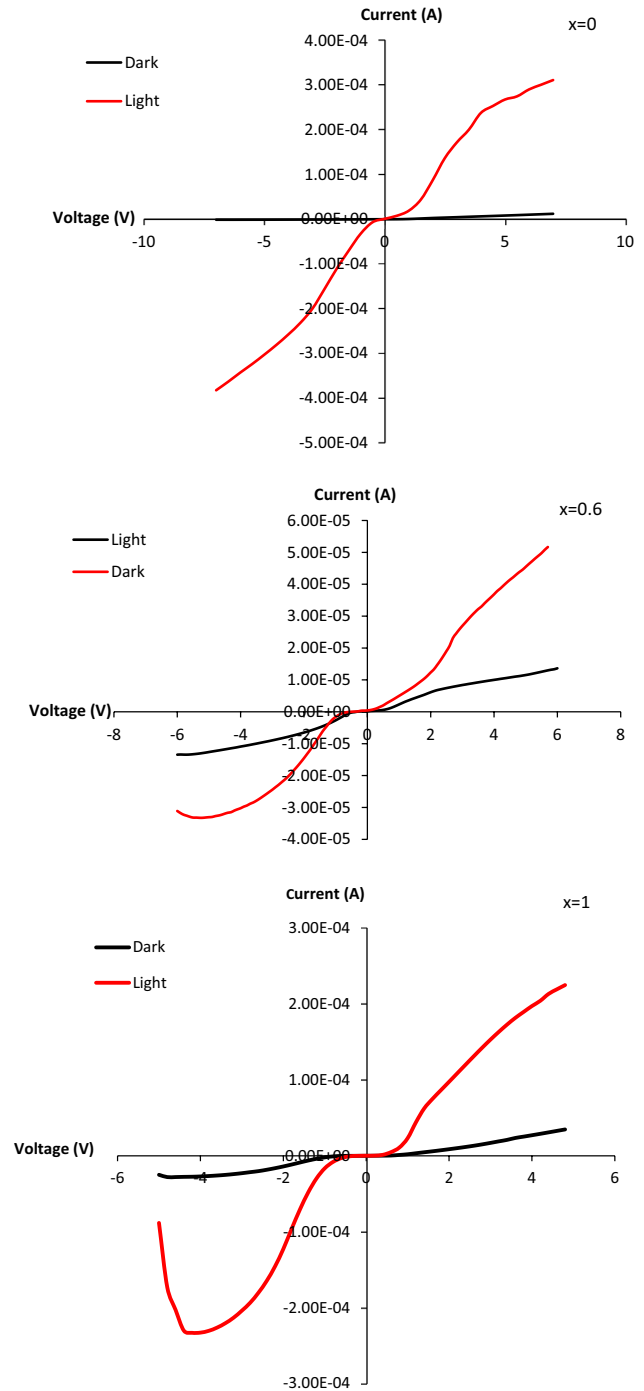
1, 2 for group IV, III–V, and II–VI semiconductors, respectively. In Table 1, the calculated bulk modulus values are investigated. We may conclude that the present bulk moduli calculated in a different way than the definition of Cohen exhibit the same chemical trends as those found for their constituents. It is observed that Cd content is proportion inversely with the stiffness of CZCTS quinary alloy nanostructures.

Figure 3 shows top-view of FE-SEM images of  $\text{Cu}_2\text{Zn}_{1-x}\text{Cd}_x\text{SnS}_4$  quinary alloy nanostructures grown on PS (63.93 %) with different Cd content ( $x$ ). At  $x = 0$ , the surface has multiple pores among the grains. The number of pores decreases as the Cd content increases. The sample grown with  $x = 0.6$  clearly to show the largest grain size. This is because of the adjacent  $\text{Cu}_2\text{Zn}_{1-x}\text{Cd}_x\text{SnS}_4$  grains are tended to merge into larger ones as Cd content increases. Contrary, at  $x = 1$ , the grains have extremely small pores. In addition, the top-view image of as prepared n-PS/ $\text{Cu}_2\text{Zn}_{1-x}\text{Cd}_x\text{SnS}_4$  reveals a uniform distribution of pores.

The PL spectra of  $\text{Cu}_2\text{Zn}_{1-x}\text{Cd}_x\text{SnS}_4$  quinary alloy nanostructure deposited on PS prepared at the current density of  $5 \text{ mA/cm}^2$  for 120 min are shown in Fig. 4. The peak position ( $\lambda_{\text{max}}$ ) and the full width at half maximum (FWHM) of the band gap are shifted with increasing Cd content from  $\lambda_{\text{max}} = 672.87 \text{ nm}$  ( $x = 0$ ) to  $677 \text{ nm}$  ( $x = 0.6$ ) to  $701 \text{ nm}$  ( $x = 1$ ) corresponding to energy gaps  $1.84 \text{ eV}$  ( $x = 0$ ),  $1.83 \text{ eV}$  ( $x = 0.6$ ) and  $1.76 \text{ eV}$  ( $x = 1$ ), respectively. The shifting is due to substituting Zn atoms with Cd atoms to produce a lower energy gap, because ZnS has a direct energy gap ( $3.62 \text{ eV}$ ) (Hu et al. 2013; Luque et al. 2013). The current-to-voltage (I–V) characteristics of the fabricated devices in dark and white light (tungsten lamp) under  $460 \text{ nm}$  for the Ag/n-PS/ $\text{Cu}_2\text{Zn}_{1-x}\text{Cd}_x\text{SnS}_4$ /Ag diode at  $x = 0, 0.6, 1$  are shown in Fig. 5. The photocurrent ratio between  $I_{\text{dark}}$  and  $I_{\text{ph}}$  is increased with increasing Cd



**Fig. 4** PL of  $\text{Cu}_2\text{Zn}_{1-x}\text{Cd}_x\text{SnS}_4$  quinary alloy nanostructures at different Cd contents,  $x = 0, 0.6, 1$



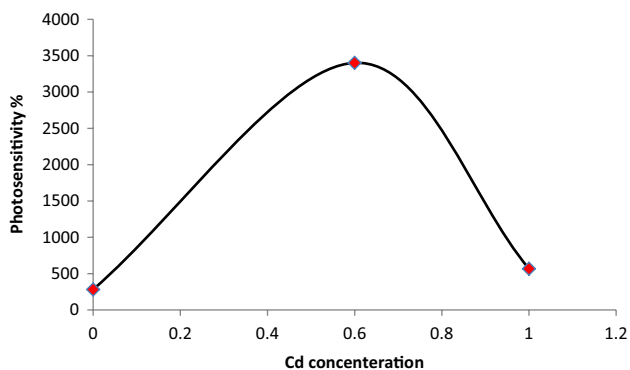
**Fig. 5** I–V characteristics of the Ag/n-PS/ $\text{Cu}_2\text{Zn}_{1-x}\text{Cd}_x\text{SnS}_4$ /Ag diode at different Cd contents,  $x = 0, 0.6, 1$ . The measurements are in dark and under white light

content. The photosensitivities of the heterojunctions were calculated by the equation

$$S = \frac{I_{\text{ph}} - I_{\text{dark}}}{I_{\text{dark}}} \times 100 \% \quad (5)$$

**Table 2** Photoresponse properties of the Ag/n-PS/Cu<sub>2</sub>Zn<sub>1-x</sub>Cd<sub>x</sub>SnS<sub>4</sub>/Ag diode 1 at bias voltage 3 V and wavelength 490 nm

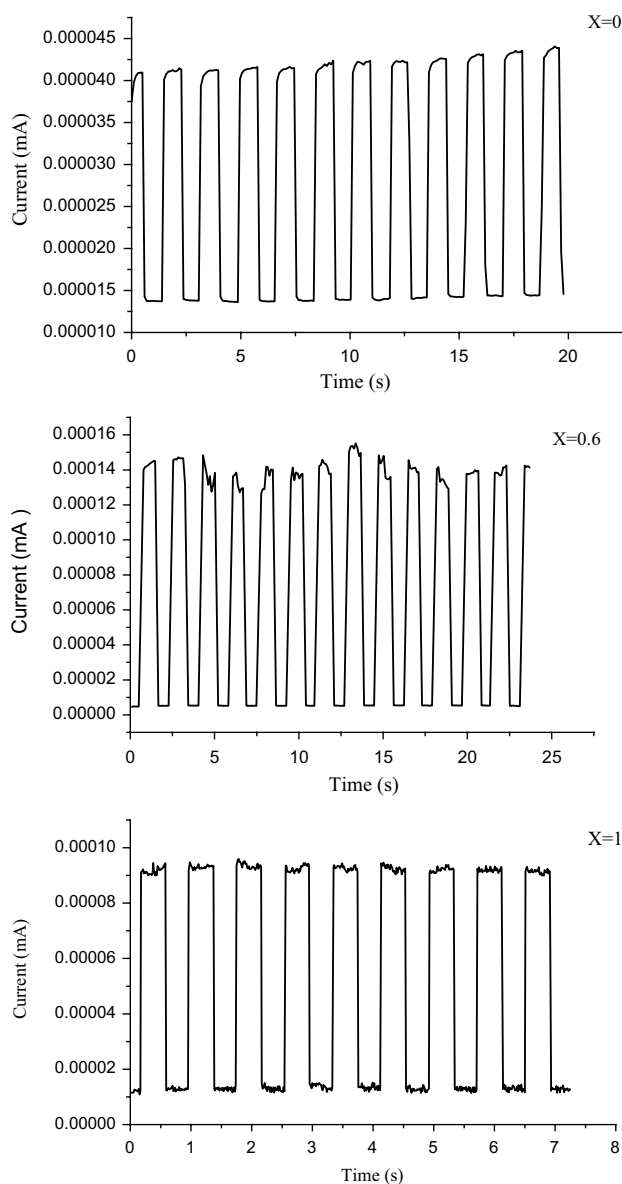
X	I <sub>dark</sub> (A)	I <sub>ph</sub> (A)	I <sub>ph</sub> /I <sub>dark</sub>	S <sub>ph</sub> (%)
0	1.08 × 10 <sup>-5</sup>	4.13 × 10 <sup>-5</sup>	3.82	282.40
0.6	7.35 × 10 <sup>-6</sup>	25 × 10 <sup>-5</sup>	34.01	3401.36
1	3.28 × 10 <sup>-5</sup>	2.19 × 10 <sup>-4</sup>	6.67	567.68



**Fig. 6** Photosensitivity o of the Ag/n-PS/Cu<sub>2</sub>Zn<sub>1-x</sub>Cd<sub>x</sub>SnS<sub>4</sub>/Ag diode at different Cd contents, x = 0 0.6, 1

Based on the I–V curve (Fig. 5) and Eq. (5), the I<sub>dark</sub> and I<sub>ph</sub> are listed in Table 2. The dark current is decreased with increasing Cd content. It can be attributed to the improvement of the crystalline properties with increasing of Cd content, that leading to the reduction of the structural defects of the surface that guide to improvement in the electrical properties. Figure 6 shows that the photosensitivity increases as Cd content increases to 3401.36 for x = 0.6 and decreases to 567.68 to x = 1 (Table 2). The photo-to-dark current ratio was achieved at maximum of 34.01 for x = 0.6 compared with 3.82 for x = 0 (Table 2). No more increase in the output current with increasing Cd content up to x = 1, this is related to the substitution impurities relating to trapping centre production which inversely affect the photocurrent. The photocurrent values of Cu<sub>2</sub>Zn<sub>0.4</sub>Cd<sub>0.6</sub>SnS<sub>4</sub> quinary alloy nanostructures are higher than their endmembers Cu<sub>2</sub>ZnSnS<sub>4</sub> and Cu<sub>2</sub>CdSnS<sub>4</sub>.

To confirm the better performance of the photodetector and reproducibility of the device (Ag/n-PS/Cu<sub>2</sub>Zn<sub>1-x</sub>Cd<sub>x</sub>SnS<sub>4</sub>/Ag), it was examined by cyclically switching the white light on and off. Figure 7 shows the I<sub>ph</sub> as a function of time intervals (I–t) of Cu<sub>2</sub>Zn<sub>1-x</sub>Cd<sub>x</sub>SnS<sub>4</sub> quinary alloy nanostructures at x = 0, 0.6, 1. I<sub>ph</sub> sharply increased/decreased to reach the maximum/minimum under 490 nm white light on/off. The Cu<sub>2</sub>Zn<sub>1-x</sub>Cd<sub>x</sub>SnS<sub>4</sub> quinary alloys nanostructures at x = 0.6 (Fig. 7b) shows better performance than at x = 0, 1 (Fig. 7a, c). The result indicated



**Fig. 7** Photocurrent response spectra of Ag/n-PS/Cu<sub>2</sub>Zn<sub>1-x</sub>Cd<sub>x</sub>SnS<sub>4</sub>/Ag diode different Cd contents, x = 0, 0.6, 1 with white illumination, 490 nm turned on and off repeatedly at bias voltage 3 V

**Table 3** Response and decay time correspond to I<sub>ph</sub>/I<sub>dark</sub> ratio of Ag/n-PS/Cu<sub>2</sub>Zn<sub>1-x</sub>Cd<sub>x</sub>SnS<sub>4</sub>/Ag diode

X	Response time (t <sub>Res</sub> ) (s)	Decay time (t <sub>Rec</sub> ) (s)	I <sub>ph</sub> /I <sub>dark</sub>
0	0.036	0.025	2.85
0.6	0.010	0.011	14
1	0.021	0.017	8.18

an improvement with increasing Cd content, x = 0.6. The I<sub>ph</sub>/I<sub>dark</sub> ratio was 2.85, 14 and 8.18 for x = 0, 0.6 and 1, respectively as given in Table 3. The calculated response

time ( $t_{Res}$ ) and decay time ( $t_{Rec}$ ) using 490 nm and 3 V show that all values are decreased with increasing of Cd content indicating an improved photoresponse. Performance of  $Cu_2Zn_{1-x}Cd_xSnS_4$  quinary alloys nanostructures shows  $t_{Res}$  (0.010) and  $t_{Rec}$  (0.011) for  $x = 0.6$  (Table 3).

#### 4 Conclusion

The  $Cu_2Zn_{1-x}Cd_xSnS_4$  quinary alloys nanostructures are grown on n-PS (63.93 %) using spin coating technique. It is deduced that Cd contents correlate directly with lattice constant and inversely with bulk modulus, that reflect the stiffness of  $Cu_2Zn_{1-x}Cd_xSnS_4$  quinary alloys nanostructures. The band gap is found to decrease with increasing Cd contents. Photocurrent measurements show the highest photoresponse for  $Cu_2Zn_{0.4}Cd_{0.6}SnS_4$  quinary alloys nanostructures.

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